Storage and Release of Soluble Hexavalent Chromium from Chromate Conversion Coatings on Al Alloys

Kinetics of Release

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The release of chromate ions from chromate conversion coatings (CCCs) on Al alloys was studied, and the effect of aging of CCCs on the chrome release kinetics was investigated. Chromate release from CCCs into aqueous solutions was measured by monitoring the change in the chromate concentration in solution using UV-visible spectroscopy. Heat-treatment of the CCC greatly reduced the chromate release rate. The chromate release rate also decreased with increasing aging time at room temperature. A diffusion-control model was proposed based on the notion that the CCC in an aqueous solution is a porous, two-phase structure consisting of a solid phase with adsorbed Cr(VI) species that is in local Langmuir-type equilibrium with an interpenetrating solution phase. This model results in a concentration gradient of soluble Cr(VI) in the solution phase of the CCC as chromate is released. The concentration and diffusion coefficients of soluble Cr(VI) in CCC were estimated. The estimated diffusion coefficient tended to decrease with aging time, suggesting that the CCC is modified with aging time.

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For aircraft coating technology, understanding the mechanism of aluminum corrosion and the effect of chromates on corrosion inhibition is important. Chromate conversion coatings (CCCs) are commonly used on aluminum alloy surfaces to improve the corrosion resistance and to establish an adhesive base for organic coatings. Chromate ions have increased the pitting potential of aluminum alloys in close proximity to a CCC-treated sample. The release of Cr(VI) from CCCs was further studied using UV-Visible spectroscopic analysis of solutions exposed to samples with CCCs. The observations were consistent with a release mechanism similar to Langmuir adsorption and desorption of Cr(VI) from a porous, insoluble Cr(III) matrix. To fully understand the inhibition behavior of CCCs, one must fully understand the chromate release phenomenon. In the present study, the release of Cr(VI) from CCCs is examined for samples under a range of conditions, and a model describing the release kinetics is developed.

Experimental

CCCs were prepared on aluminum alloys AA1100 and AA2024-T3 by immersion in Alodine 1200S (Henkel Surface Technologies) solutions at room temperature. The Alodine bath consisted of 9.0 g L-1 of a mixture reported to be 50-60% CrO3, 20-30% KBF4, 10-15% K2Fe(CN)6, 5-10% K2ZrF6, and 5-10% NaF by weight, which is approximately 0.065 M in Cr. The pH of the bath was adjusted to 1.5 by the addition of nitric acid. The AA1100 samples (3 cm × 8.35 cm × 50 μm) were given an acidic pretreatment using a triply diluted, commercially available detergent, Alumiprep 33 (Henkel Surface Technologies), and the AA2024-T3 samples were prepared by polishing to 1200 grit paper in water, rinsing with deionized (DI) water and degreasing in alcohol before the Alodine treatment. Most samples were treated in the Alodine bath for 5 min. After Alodine treatment, the samples were rinsed by dipping in DI water, and aged in air at room temperature or at 50, 100, or 200°C. After a certain period of aging time, chromate release was determined during immersion in stagnant DI water or 0.1 M NaCl solution at room temperature in glass beakers or glass vials. To avoid evaporation of water, the beakers were covered with vinyl sheets.

The concentrations of released Cr ions in solutions were periodically measured by UV-visible spectroscopy using a home-built system controlled by a LabVIEW program. A cuvette of 10 mm path-length containing sample or reference solution was placed in the path of monochromated incident light. The intensity of the transmitted light was measured by a power meter, and the absorption was determined as a function of the wavelength of the incident light. After the measurements, the solution in the cuvette was replaced in the same beaker. Quantitative determination of Cr(VI) concentration was calculated from UV absorption at the isosbestic point of 340 nm, where the absorption is independent of pH and is linearly related to the concentration of total Cr(VI). The linear calibration relation of Cr(VI) concentration to absorption measured at 340 nm can be written as [Cr(VI)] = (Abs340/0.00157)/1062. An absorption of 1 is equal to 1 mM of total Cr(VI); hereafter, we express the concentration of total Cr(VI) by absorption.

Results and Discussion

Effect of heat-treatment.—The 0.1 M NaCl solution in contact with a CCC-treated AA1100 sample aged for 2 h at room temperature tended to show an increase in Cr(VI) concentration with time, approaching a constant value, as evidenced by an increase in the absorption at 340 nm (see Fig. 1). The sample heat-treated at 50°C for 2 h showed remarkably less chrome release than the sample aged at room temperature. CCC samples heat-treated at 100 or 200°C did not release detectable amounts of chromates. The CCC sample aged at 200°C lost the brownish golden color characteristic of CCCs as the color of the sample turned to metallic gray. Clearly, heat-treatment suppressed the release of chromates from the CCCs.

Effect of aging and ratio of surface area to solution volume (S/V).—Figure 2 shows the effect of room-temperature aging on chromate release from CCC-treated AA1100. Note that all samples described hereafter were aged at room temperature. In general, the concentration of Cr(VI) in solution increased with immersion time
and then reached an almost constant level of concentration. The release rate of chromates from CCC decreased with increasing room-temperature aging time. The release rate in 0.1 M NaCl solution was higher than that in DI water. The samples polished slightly by SiC paper before immersion in water did not behave significantly differently than the unpolished samples, which indicates that aging did not only affect the very surface of the CCCs. For the samples in this figure, the surface area and solution volume were 50 cm$^2$ and 30 mL, respectively, so that $S/V$ was 1.67 cm$^{-1}$.

The release of Cr(VI) from CCCs on AA1100 aged for 2 h for different $S/V$ ratios is compared in Fig. 3. The Cr(VI) concentration increased with $S/V$ ratio. The steady-state value of Cr(VI) concentration for $S/V = 5$ was about three times the value for $S/V = 1.67$, so it varied linearly with this ratio. However, this linear dependence was not maintained for higher $S/V$ ratios. The $S/V$ ratios of 10 and 20 cm$^{-1}$ were obtained by coiling the strip sample and immersing it in a glass vial with a relatively little amount of solution. Because of the coiling, it is possible that not all regions of the sample surface had access to the water or that the Cr(VI) concentration was inhomogeneous in the solution. Nonetheless, the measured concentration of Cr(VI) continued to increase with $S/V$ and did not exhibit a limiting value indicative of solution saturation. Accordingly, the plateau of the chromate release curves with time must be determined by something other than saturation of the solution with chromate. One possibility, as described the section on the release model, is that the chromate solution concentration limit is a result of a Langmuir-type adsorption equilibrium with chromate in the CCC.$^4$

Careful examination shows that the concentration of Cr(VI) in Fig. 3 for $S/V = 1.67$ cm$^{-1}$ is roughly twice that in Fig. 2 for the same $S/V$. The samples in Fig. 2 were treated in an Alodine solution stored for a long period of time, about 1 year, whereas the samples in Fig. 3 were made in fresh Alodine solutions. The details of sample preparation obviously can greatly affect their behavior. We were unable to investigate this discrepancy further; the samples discussed hereafter were all made in fresh Alodine solutions.

**Effect of substrate.**—Figure 4 shows Cr(VI) release from CCCs formed on AA1100 and AA2024-T3 substrates. The release of
chromates from CCCs on AA2024-T3 is slightly lower than that from CCCs on AA1100, which indicates that the CCCs on different substrates are different in nature. McCreery and co-workers have shown that the accelerating influence of ferricyanide in Alodine solutions is reduced on Cu and Cu-rich surfaces, such as would be found at the intermetallic particles, as a result of adsorption of ferricyanide at the surface, thus decreasing Cr(III) generation and CCC formation. The CCC on AA2024-T3 prepared by b5 mm immersion in Alodine bath showed a higher release of chromate than that from a sample given only a 1 min immersion. This suggests that the amount of Cr(VI) increased with increasing immersion time in the Alodine solution.

Release kinetics.—To see the effect of the S/V ratio on release behavior, the UV absorption was normalized. Figure 5 shows a comparison of the chromate release for S/V = 5 cm⁻¹ and three times the chromate release for S/V = 1.67 cm⁻¹. After this normalization, the release curves were almost identical. Therefore, the kinetics of chromate release and the total amount of chromate (not the concentration in solution) were approximately independent of S/V.

To further investigate the cause of the plateau in the release curves, a second sample made by the same procedure was added to a cell containing a sample that had been exposed for about 400 h of immersion at S/V = 5 cm⁻¹ (see Fig. 6). The first sample was kept in the solution, so that the final S/V ratio was 10 cm⁻¹. The concentration of Cr(VI) started to increase again after the addition of the second sample. In another experiment, the solution of a sample with S/V = 1.67 cm⁻¹ was exchanged for fresh DI water after more than 300 h of immersion. A very small amount of Cr(VI) was found in the solution after the exchange. The further increase in chromate concentration with the addition of another sample supports our previous suggestion that the plateau concentration of Cr(VI) does not correspond to saturation. The fact that additional chromate was released when an equilibrated sample was placed in a fresh water solution indicates that the original steady state was not determined by depletion of chromate in the CCC. However, the amount released into the replaced DI water was less than that released into the original solution because the Cr(VI) concentration in the film was lower as a result of the first exposure.

The release curves were normalized in another fashion in Fig. 7. Here, the concentration was normalized by the concentration at the plateau, C(∞), to study the release kinetics. Curves of normalized concentration changed considerably with aging time, but they were not strongly influenced by the S/V ratio. This suggests that the release of Cr(VI) was related to the structure of the film, which may have affected diffusion of chromates in the CCC.

Release model.—In a previous paper, the steady-state solution concentrations were shown to be a result of a Langmuir-type equilibrium with Cr(VI) species adsorbed on the CCC. However, the fact that a sample treated for a longer time resulted in a higher solution concentration suggests that the CCC is not a two-dimensional surface with chromate adsorption sites, but that Cr(VI) binds to Cr(III) throughout a three-dimensional (3D) porous Cr(III) film. A diffusion model was developed assuming Langmuir-type equilibrium between dissolved Cr(VI) and a 3D CCC film. This model accounts for all of the observations previously described.

The basis for the model is shown in Fig. 8. The structure of a CCC in contact with an aqueous solution is considered to be a 3D porous Cr(III) matrix or backbone with chromate adsorption sites and an interpenetrating aqueous phase. The chromate concentrations in the aqueous phase within the CCC and the adsorbed Cr(VI), [Cr(VI)(aq)] and [Cr(VI)(ad)], respectively, can be functions of time and depth from the metal/CCC to the CCC/bulk-solution interfaces. The depth distribution of Cr(VI) is assumed to change with immersion time, as shown by the curves in the figure. Cr(VI) at the CCC/bulk-solution interface dissolves into the bulk solution and is replenished.
It is assumed that this Langmuir-type adsorption/desorption equilibrium is fast compared to diffusion in the aqueous phase of the CCC. The binding energy of the adsorbed species is large corresponding to the formation of covalent bonds.\(^4\) This means that the rate constant for adsorption is much higher than that of desorption, so \([\text{Cr(VI)} (\text{aq})] < [\text{Cr(VI)} (\text{ad})]\), and the effective diffusivity should be relatively small.

\([\text{Cr(VI)} (\text{aq})]\) and \([\text{Cr(VI)} (\text{ad})]\) are related locally in the CCC according to the equilibrium of Reaction 1. Because of the interchange of Cr(VI) ions between the adsorption and desorption state, the diffusion equation for the system of chromate release from CCC is more complicated than the usual type of Fickian diffusion.

As discussed in detail in the Appendix, the diffusion of Cr(VI) species in CCC is described by

\[
\frac{\partial C_f}{\partial t} \left[1 + \frac{\eta \rho}{(C_f + \eta \rho)^2}\right] = D \frac{\partial^2 C_f}{\partial x^2} \tag{2}
\]

where \(C_f (=[\text{Cr(VI)} (\text{aq})])\) is the chromate concentration in the aqueous phase of the CCC film, \(\eta\) is the number of the adsorption sites in a unit volume, and \(\rho\) is the ratio of desorption to adsorption rate constants, \(k_2/k_1\). The concentration of adsorbed chromate, \(C_a (=[\text{Cr(VI)} (\text{ad})])\), and the total chromate concentration, \(C (=[\text{Cr(VI)} (\text{aq})] + [\text{Cr(VI)} (\text{ad})])\), at a particular \(x\) and \(t\) are known by

\[
C_a = C - C_f = \eta \theta \frac{\eta k_1 C_f}{k_1 C_f + k_2 \eta} = \frac{\eta C_f}{C_f + \eta \rho} \tag{3}
\]

The boundary conditions are as follows

\[
\frac{\partial C_f}{\partial x} (0,t) = 0, \quad C_f(L,t) = C_s(t) = -D S \int_0^t \frac{\partial C_f}{\partial x}(L,t) dt \tag{4}
\]

where \(C_s\) is the concentration of Cr(VI) in the bulk solution, \(S\) is the surface area of the CCC, \(V\) is the volume of the bulk solution, and \(L\) is the film thickness of the CCC. Due to the time-dependent boundary condition on the right side of Eq. 4, it is difficult to obtain a solution expressed by an equation and to estimate unknown parameters by curve fitting. A program was written to solve these equations and to plot the predicted solution concentration of Cr(VI) as a function of time.

It was previously reported that the chromate release was consistent with a mechanism similar to a Langmuir adsorption isotherm

\[
\Gamma_{\text{VI}} \frac{\Gamma_s - \Gamma_{\text{VI}}}{\Gamma_s} = \beta [\text{Cr}^{3+}] [\text{H}^+] = \frac{\theta}{1 - \theta} \tag{5}
\]

where \(\Gamma_{\text{VI}}\) is the coverage of Cr(VI) on the Cr(III) oxide in mol cm\(^{-2}\), \(\beta\) is a binding constant, and \(\theta\) is the fractional coverage of Cr(VI). However, when the rates of adsorption and desorption are the same

\[
k_1 [\text{Cr}^{6+}] (1 - \theta) = k_2 \eta \theta \tag{6}
\]

Note that \(k_1\) is proportional to \(\eta\). Owing to its high mobility, it is assumed that \([\text{H}^+]\) is independent of the depth and time. From Eq. 5 and 6

\[
\beta [\text{H}^+] = \frac{k_1}{k_2} \frac{1}{\eta} = \frac{1}{\rho \eta} \tag{7}
\]

It has been reported that \(\beta [\text{H}^+] = 10^3\) mol\(^{-1}\) L\(^{-1}\). Therefore, it is predicted that \(\rho \eta\) is roughly \(10^{-6}\) mol cm\(^{-2}\). Fitting of the model to the experimental data was performed assuming that \(\rho \eta\) is about \(10^{-6}\) mol cm\(^{-2}\), and the diffusion coefficient, \(D\), the initial concentration of \([\text{Cr}^{6+}]\), and the film thickness were estimated. The values of \(\rho\) (unitless) and \(\eta\) (mol cm\(^{-2}\)) were varied from \(1 \times 10^{-7}\) to \(2 \times 10^{-3}\). The unknown parameters are the diffusion coefficient, \(D\), the CCC film thickness, \(L\), the number of the adsorption sites, \(\eta\), and the ratio of rate constants, \(k_2/k_1\) (=\(\rho\)). Sets of values of these parameters were input to the program for obtaining a numerical solution of the differential equation, Eq. 4, and the predicted Cr(VI) concentrations in the bulk solution as a function of time were fitted to the experimental data, providing estimated values of the parameters. \(L\), \(\rho\), and \(\eta\) affect the plateau Cr(VI) concentration, but \(D\) does not.

Figure 9 shows experimental data and curves calculated using a representative set of estimated values of parameters. The model predictions exhibit good agreement with the experimental data, which
The effect of aging is addressed later in this section.

indicates that the proposed diffusion model and values of parameters estimated are appropriate. This supports the assumptions that the CCC is a 3D porous Cr(III) matrix or backbone with chromate adsorption sites and an interpenetrating aqueous phase, and that the diffusion coefficient decreases with the effective diffusivity within a CCC is reduced relative to that in a bulk solution by two factors: the effect of adsorption on the Cr(III) matrix, which is accounted for by the model, and the physical barrier nature of the CCC.

The difference between and typical diffusivities of ions in bulk solutions can be explained by the effect of the network structure of the Cr(III) matrix acting as a diffusion barrier. The CCC can be viewed as a biphasic structure consisting of interpenetrating solid Cr(III) matrix and aqueous channels, and the pathways for Cr(VI) diffusion are the aqueous channels. These channels are apparently so tight and convoluted that they impede the transport of the soluble Cr(VI) species through the CCC. This effect is probably larger when the local pH is low and the dominant Cr(VI) species is the larger dichromate anion. Formation of hydrogen bonds may also affect the diffusion of Cr(VI) species. Room-temperature aging probably decreases the Cr(VI) diffusion coefficient in the CCC because of the formation of a multinuclear Cr(III)O-Cr(III) structure that impedes the diffusion of Cr(VI) species. The structure of CCCs and the effects of aging have recently been studied using extended X-ray absorption fine structure analysis. Aging resulted in a decrease in the effective diffusivity in the CCC because of the extent of cross-linking increases with time or heat-treatment. It may be more appropriate to compare the diffusivity values to those for ions in polymers, which vary over a wide range, from 10⁻¹⁰ to 10⁻³ cm²/s⁻¹ in polymer electrolytes to near zero for immobile ions or neutrals. For example, polyethylene oxide electrolytes chosen for high ion mobility exhibit diffusion coefficients for small molecules of 3 × 10⁻⁷ to 2 × 10⁻⁸ cm²/s⁻¹. Much smaller diffusion coefficients are expected for cross-linked polymers such as Cr(III)(OH)₃, and the values are likely to decrease as the extent of cross-linking increases with time or heat-treatment. Therefore, the values shown in Fig. 11 are within an acceptable range for cross-linked polymers.

Figure 12 shows the effective diffusivity of Cr(VI) in the CCC on AA1100 as a function of room-temperature aging time, position within the CCC, and time of immersion in water. The effective diffusivity decreases with immersion time owing to a decrease in both [Cr(VI)(aq)] and [Cr(VI)(ad)] as the Cr(VI) is released from the CCC into the bulk solution. The effective diffusivity is lower at the CCC/solution interface than at the Al/CCC interface because of the Cr(VI) concentration gradient within the CCC, as is shown later in this section. The effect of room-temperature aging time is reflected in the change in D shown in Fig. 11.

The estimated initial concentration of soluble Cr(VI) is shown in Fig. 13. The amount of soluble chromate in the CCC does not change markedly with aging time at room temperature, up to 4 days.

\[
D_{\text{eff}} = \frac{D}{1 + \left( \frac{\eta}{\eta_v} \right)^2}
\]

\(D_{\text{eff}}\) is different than \(D\) because of the interaction of Cr(VI) with the Cr(III) matrix; \(D\) is the diffusion coefficient of Cr(VI) species in the aqueous phase of CCCs with no consideration of the influence of adsorption of Cr(VI) with the Cr(III) matrix. Figure 11 shows that the values for \(D\) are on the order of 10⁻¹¹ cm²/s⁻¹, which is significantly lower than that expected for the diffusivity of Cr(VI) ions in bulk solutions. Therefore, the effective diffusivity within a CCC is reduced relative to that in a bulk solution by two factors: the effect of adsorption on the Cr(III) matrix, which is accounted for by the model, and the physical barrier nature of the CCC.

The difference between \(D\) and typical diffusivities of ions in bulk solutions can be explained by the effect of the network structure of the Cr(III) matrix acting as a diffusion barrier. The CCC can be viewed as a biphasic structure consisting of interpenetrating solid Cr(III) matrix and aqueous channels, and the pathways for Cr(VI) diffusion are the aqueous channels. These channels are apparently so tight and convoluted that they impede the transport of the soluble Cr(VI) species through the CCC. This effect is probably larger when the local pH is low and the dominant Cr(VI) species is the larger dichromate anion. Formation of hydrogen bonds may also affect the diffusion of Cr(VI) species. Room-temperature aging probably decreases the Cr(VI) diffusion coefficient in the CCC because of the formation of a multinuclear Cr(III)O-Cr(III) structure that impedes the diffusion of Cr(VI) species. The structure of CCCs and the effects of aging have recently been studied using extended X-ray absorption fine structure analysis. Aging resulted in a decrease in the effective diffusivity in the CCC because of the extent of cross-linking increases with time or heat-treatment. Therefore, the values shown in Fig. 11 are within an acceptable range for cross-linked polymers.

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The estimated initial concentration of soluble Cr(VI) is shown in Fig. 13. The amount of soluble chromate in the CCC does not change markedly with aging time at room temperature, up to 4 days.
This suggests that the Cr(VI) is not reduced to Cr(III) during room-temperature aging, does not form an insoluble product, and does not form an irreversible bond with the Cr(III) matrix species. Note that the situation changes after very long aging times. A CCC-treated panel of AA3003-H14 contained virtually no Cr(VI) by X-ray absorption near-edge structure analysis.\textsuperscript{15}

The estimated thickness of the CCC is shown in Fig. 14 as a function of aging time. The estimated CCC thickness was almost constant at approximately 1 \( \mu \text{m} \), which corresponds with the reported CCC thickness.\textsuperscript{5}

Figure 15 shows the change in the concentration of adsorbed Cr(VI) on adsorption sites of CCC and free Cr(VI) in aqueous phase in CCC as a function of immersion time. Because the binding of Cr(VI) to Cr(III) is strong,\textsuperscript{4} Cr(VI) species exist primarily as adsorbed Cr(VI) in the CCC. The shape of the adsorbed Cr(VI) curves is similar to those of free Cr(VI) because local equilibrium is assumed. The relationship between adsorbed Cr(VI) and free Cr(VI) is expressed by Eq. A-7 in the Appendix. The depth profiles of free Cr(VI) in CCC aged for 2 or 96 h derived from the model calculation are shown in Fig. 16 (top) and (bottom), respectively. Because the main difference between 2 and 96 h aged CCCs is in the value of the diffusion coefficient, the shapes of the depth profiles are basically similar to each other, although the time scale is different. The Cr(VI) concentration at the surface of the CCC at the initial stages of immersion is smaller than that at longer times because the Cr(VI) concentration in bulk solution is initially 0 and gradually increases by the Cr(VI) release from CCC. At higher \( S/V \) ratios, the change of Cr(VI) concentration at the CCC surface is larger.

\textit{Implications of model.}—This release model supports the view of a CCC as a 3D porous Cr(III) matrix or backbone with chromate adsorption sites and an interpenetrating aqueous phase. Release of Cr(VI) from this biphase structure is driven by a Langmuir-type equilibrium of the CCC at the film/bulk-solution interface, and the release kinetics are determined by the diffusion of chromate in the aqueous phase in conjunction with a local Langmuir-type equilibrium.

From the corrosion point of view, this model provides important implications as to how Cr(VI) releases in the field. SrCrO\textsubscript{4} pigment in paint does not form a Langmuir-type interaction with the solution in which it is contact; chromate pigment releases Cr(VI) until the solution is saturated. When Cr(VI) in the surface solution reduces to Cr(III) for corrosion protection of a nearby metal, the remaining SrCrO\textsubscript{4} solid releases more Cr(VI) to maintain saturation. The solution remains saturated until the SrCrO\textsubscript{4} is fully consumed. For

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{fig12}
\caption{Effective diffusion coefficient of Cr(VI) species in the CCCs on AA1100 as a function of aging time at room temperature, location within the CCC, and time of immersion in water.}
\end{figure}

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{fig14}
\caption{Estimated thickness of CCCs on AA1100 aged for 2, 24, or 96 h.}
\end{figure}

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{fig15}
\caption{Change in the concentration of adsorbed Cr(VI) and free Cr(VI) in aqueous phase at the Al1100/CCC interface and at the CCC/solution interface for CCCs aged 2, 24, or 96 h. The curves were obtained by using the values of parameters used in Fig. 9.}
\end{figure}
chromate in pigmented paint is usually far in excess of the amount of Cr release. As Cr~VI~CCCs, the equilibrium Cr~VI~ solution concentration depends on the S/V ratio, surface Cr~VI~ concentration, and solution pH, etc. Release of Cr~VI~ from the CCC film results in a lower surface Cr~VI~ concentration. As Cr~VI~ in the solution is consumed, the equilibrium Cr~VI~ solution concentration decreases. The total amount of chromate in pigmented paint is usually far in excess of the amount in a CCC. However, for a CCC and SrCrO$_4$ pigment containing the same amount of Cr~VI~, the CCC film lasts longer because of the different type of equilibrium reached with the solution.

As shown in Fig. 2, NaCl addition to DI water enhanced Cr~VI~ release. It seems that chloride anions inhibit the reversible binding of Cr~VI~ to hydrated chromium hydroxide, so that the diffusivity of chromates is enhanced. The higher release rate in the presence of corrosive chloride ions is useful from the viewpoint of corrosion protection.

**Conclusions**

Release of chromate from CCC was verified by UV absorption measurements. Heat-treatment of CCC markedly reduced the ion release rate. The release rate of chromates from CCC decreased with aging time at room temperature.

A diffusion model was proposed in conjunction with a local Langmuir-type equilibrium with an interpenetrating solution phase. Using this model, the diffusion coefficient of chromates in the CCC, the CCC thickness, and the initial concentration of Cr~VI~ were estimated. The diffusion coefficient decreased with aging time at room temperature, presumably because of the development of the CCC matrix.

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**Appendix**

**Diffusion of Hexavalent Chromium in CCC**

The total concentration of Cr~VI~ within the coating is a function of the position, $x$, and the immersion time, $t$, and is given by $C(x,t)$, with $x = 0$ at the metal/coating interface and $x = L$ at the coating/solution interface. In general, $C$ is comprised of two contributions, i.e., the concentration $C_f$ of Cr~VI~ that is free to diffuse, and the concentration $C_a$ that is adsorbed on any of the $n_a$ adsorption sites available per unit volume of the film. We assume that $n_a$ is independent of $x$ and $t$. Hence

$$C(x,t) = C_f(x,t) + C_a(x,t)$$  \[A-1\]

**Adsorption kinetics**—We assume that $C_f$ and $C_a$ are related by simple Langmuir-type adsorption kinetics. In that regard, consider a thin “slab” of coating extending from $x$ to $x + dx$. The number of adsorption sites per unit area of this slab is $\eta dx$. Let $\theta(x,t)$ be the fraction of these sites occupied by adsorbed Cr~VI~, from which we deduce that

$$C_a dx = \eta dx \theta \text{ or } C_a = \eta \theta$$  \[A-2\]

We consider the insoluble Cr~III~ to have many surface hydroxyl groups available for formation of Cr~III~-O—Cr~III~ bonds. These hydroxyl groups act as sites for adsorption of Cr~VI~, accompanied by the addition of a hydrogen ion to produce water. The reaction may be represented by

$$\text{Cr}^{III} - \text{OH} + \text{Cr}^{VI}(aq) \xrightarrow{k_1} \text{Cr}^{III} - \text{O} - \text{Cr}^{VI}(ad) + \text{H}_2\text{O}$$  \[A-3-1\]

Cr~III~—OH represents insoluble Cr~III~ oxyhydroxide, Cr~VI~(aq) is solution-phase Cr~VI~ species which include oxygen atoms, and Cr~III~—O—Cr~III~ is the adsorbed state of Cr~VI~ to Cr~III~ by oxygen. Note that Cr~VI~(aq) has one oxygen atom more than Cr~VI~(ad).

The reaction is reversible, so the CCC may release Cr~VI~, and the Cr~III~ oxyhydroxide may adsorb Cr~VI~ from the aqueous phase. Considering the existence of Cr$_2$O$_7^{2-}$ as a Cr~VI~ species, the reaction of insoluble Cr~III~ oxyhydroxide with Cr$_2$O$_7^{2-}$ can be expressed by Reaction A-3-2. Two Cr~III~ are connected with each other by two o-bridges

$$2\text{Cr}^{III} - \text{OH} + \text{Cr}^{VI}(O_7^{2-})$$

$$+ 2\text{H}^+(aq) \xrightarrow{k_1} 2[\text{Cr}^{III} - \text{O} - \text{Cr}^{VI}(ad)] + 2\text{H}_2\text{O}$$  \[A-3-2\]

From the stoichiometric point of view, Reactions A-3-1 and A-3-2 are basically the same.

The adsorption/desorption reactions within the slab are described as follows

Rate of adsorption = $k_1 C_f(1 - \theta)$  \[A-4\]

Rate of desorption = $k_2 C_a = k_2 n_a \theta$  \[A-5\]

where $k_1$ and $k_2$ are rate constants, with $k_1$ being proportional to [H$^+$]. It is assumed that $C_f > C_a$ because the binding energy of the adsorption is high, corresponding to chemisorption. We assume that [H$^+$] is independent of $x$ and $t$. Langmuir-type kinetics involves taking the two rates to be equal

$$0 = \frac{k_1 C_f}{k_1 C_f + k_2 n_a} \quad \text{or} \quad \frac{k_1 C_f}{k_2 n_a} = \frac{0}{1 - \theta}$$  \[A-6\]

We can now establish a relationship between $C_f$ and $C_a$, i.e.

$$C_a = C - C_f = \eta \theta = \frac{\eta k_1 C_f}{k_2 n_a + k_2 n_a} = \frac{\eta k_1 C_f}{k_2 n_a + \eta \theta}$$  \[A-7\]

where $\rho = k_2/k_1$. Note that $C_a < \eta$, as expected. We now have the following quadratic equation

$$C_a = C - C_f = \eta \theta = \frac{\eta k_1 C_f}{k_2 n_a + k_2 n_a} = \frac{\eta k_1 C_f}{k_2 n_a + \eta \theta}$$  \[A-7\]
\[(C - C_i)(C_j + \eta \rho) = \eta C_i \]

or \[C_i^2 + C_i(\eta \rho + \eta - C) - \eta \rho C = 0 \]  
only one solution of which is physically acceptable, that being
\[C_i = \frac{1}{2} \left( C - \eta \rho + \sqrt{(\eta + \eta \rho - C)^2 + 4 \eta \rho C} \right) \]  

[A-8]

The other solution would give \(C_i < 0\). Thus, we have \(C_i > 0\), and we can show from Eq. A-7 that \(C_i < C\), leading to the conclusion that \(0 < C_i < C\), all of which are physically correct.

The corresponding relationship for \(C_i\) is easily obtained, i.e.
\[C_i = C - C_i - \frac{1}{2} (C + \eta \rho - \sqrt{(\eta + \eta \rho - C)^2 + 4 \eta \rho C}) \]  

[A-9]

The diffusion equation—The diffusion equation appropriate to this system is understandably more complex than the “usual” Fickian-type expression, due to the interchange of \(C_{VI}\) between the free (diffusing) and adsorbed (bound) states. However, it can be derived based on a simple conservation-of-matter argument. Toward that end, consider again the slab of film extending from \(x\) to \(x + dx\), taking the surface area of the film to be \(S\). Then
\[\frac{\partial C}{\partial t}(x,t)dx = \text{Rate at which } C_{VI} \text{ is being added to the slab} \]  

[A-10]

The only way this addition of matter can occur is by the diffusive flow of free \(C_{VI}\), so that
\[\frac{\partial C}{\partial t}(x,t)dx = -D \left( \frac{\partial C}{\partial x}(x,t)S \right) \]  

where \(D\) is the diffusivity of the free \(C_{VI}\). The first term on the right side of Eq. A-12 is the rate of influx of free \(C_{VI}\) into the slab, and the second term is the rate of egress. Expanding that second term in a Taylor series and combining terms
\[\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2} \]  

[A-11]

dropping the \((x,t)\) arguments. The second term of Eq. A-13 is recognizable as Fickian diffusion in the presence of a sink. Dependence upon \(C_i\) in Eq. A-13 can be eliminated using Eq. A-7, which relates \(C_i\) to \(C_i\). Thus we obtain
\[\frac{\partial C_i}{\partial t} = D \frac{\partial^2 C_i}{\partial x^2} \]  

[A-12]

Knowing \(C_i\) at a particular \(x\) and \(z\) also tells us the values of \(C_i\) and \(C_i\) at that position and time via Eq. A-7.

Equation A-14 can be scaled, and in so doing, all explicit dependence on the parameters \(\eta\) and \(D\) is eliminated. Toward this end, we define the following dimensionless variables
\[\xi = \eta^{1/3}x, \quad \tau = \eta^{2/3}Dt, \quad \chi_i = \eta^{-1}C_i \]  

[A-13]

in terms of which Eq. A-14 is reduced to
\[\frac{\partial \chi_i}{\partial \xi} = \frac{\partial^2 \chi_i}{\partial \xi^2} \]  

[A-14]

Some of our earlier results can also be scaled if we first define
\[\chi_i = \eta^{1/3}C_i, \quad \chi = \eta^{-1}C \]  

[A-15]

in which case Eq. A-4, A-9, and A-10 become, respectively
\[\theta = \frac{k_{1/3}X_{1/3}}{k_{1/3}X_{1/3} + k_{2/3}} \quad \text{or} \quad \frac{k_{1/3}X_{1/3}}{k_{1/3}X_{1/3} + k_{2/3}} = \frac{\theta}{1 - \theta} \]  

\[\chi_i = \frac{1}{2} \left( x - 1 - \rho + \{1 + \rho - (1 + \rho - x)^2\}^{1/2} \right) \]  

[A-16]

\[\chi = \frac{1}{2} \left( x + 1 + \rho - \{1 + \rho - x)^2 + 4 \rho \chi \}^{1/2} \right) \]  

[A-17]

[A-18]

Initial and boundary conditions.—The following initial conditions are imposed
\[C(x,0) = C_i, \quad C_i(0) = 0 \]  

[A-19]

where \(C(t)\) is the concentration of \(C_{VI}\) in the solution beyond the film, which occupies volume \(V\). Moreover, it is assumed that \(C_i(x,0)\) and \(C_i(x,0)\) are related to \(C(x,0)\) through Eq. A-9 and A-10, respectively.

The boundary conditions are as follows
\[\frac{\partial C_i}{\partial x}(0,t) = 0, \quad C_i(L,t) = C_i(t) = -D \int_0^t \frac{\partial C_i}{\partial x}(L,t)dt \]  

[A-20]

The first of the two boundary conditions is a statement of the assumption that \(C_{VI}\) does not diffuse backward, into the metal. The second is a statement of the fact that \(C_i(t)\) results from the diffusive flow of \(C_{VI}\) from the film into the solution.

The boundary conditions can be scaled to make them consistent with Eq. A-16. Toward that end, we define the following new dimensionless parameters
\[\chi_i = \eta^{-1}C_i, \quad \sigma = \eta^{2/3}S, \quad \rho = \eta V, \quad \lambda = \eta^{1/3}L \]  

[A-21]

Combining Eq. A-15 and A-23 with A-22, we find that the scaled boundary conditions are
\[\frac{\partial \chi_i}{\partial \xi}(0,\tau) = 0, \quad \chi_i(L,\tau) = \chi(L,\tau) = \frac{\sigma}{\rho} \int_0^\tau \frac{\partial \chi_i}{\partial \xi}(L,\tau)\,d\tau \]  

[A-22]

Note that there is no explicit dependence on \(D\) or \(\eta\) in Eq. A-24.

Global conservation of matter.—The diffusion equation is an expression of conservation of matter at the local or macroscopic level. A global or macroscopic expression for the conservation of matter is
\[C_{L,J} = \int_0^L C(x,t)dx + \frac{V}{S} C_i(t) \]  

[A-23]

This expression can also be presented in scaled form, first defining
\[\chi_0 = \eta^{-1}C_0 \]  

[A-24]

to obtain
\[\chi_0 = \int_0^\lambda \chi(\xi,\tau)\,d\xi + \frac{\rho}{\sigma \lambda} \chi_0(\tau) \]  

[A-25]

Rearranging Eq. A-27 and utilizing the second expression in Eq. A-24 yields
\[\chi(\lambda,\tau) = \sigma \chi_0 - \lambda \int_0^\tau \chi(\xi,\tau)\,d\xi + \chi_0(\xi,\tau)\,d\xi \]  

[A-26]

But \(\chi_0 = \chi_i(\xi + \rho)\), so differentiation with respect to \(\tau\) yields
\[\frac{\partial \chi}{\partial \xi}(\lambda,\tau) = -\frac{\sigma}{\rho} \int_0^\tau \frac{\partial \chi_0}{\partial \xi}(\xi,\tau)\left[1 + \frac{\rho}{\sigma \lambda} \chi_0(\tau)\right] d\xi \]  

[A-27]

Numerical representation.—Equation A-16 represents a partial differential equation with boundary conditions expressed by the first expression in Eq. A-24 and by Eq. A-29. This partial differential equation can be reduced to a system of ordinary differential equations by discretizing the \(\xi\) axis and assigning an ordinary differential equation to each of the points considered. To do this, divide the interval from 0 to \(\lambda\) into \(2n\) equal subintervals. Then there will be \(2n + 1\) nodes and \(2n + 1\) ordinary differential equations. Let
\[\Delta \xi = \frac{\lambda}{2n} \]  

[A-28]

and
\[\xi_i = i \Delta \xi \quad \text{for} \quad i = 0, 1, \ldots, 2n \]  

[A-29]

Let \(\chi_i, \chi_i, \chi_i\), be the values of \(\chi, \chi, \chi\), respectively, at \(\xi_i\), for \(i = 0, 1, \ldots, 2n\). The quantities \(\chi_i, \chi_i, \chi_i\) are functions of \(\tau\). In terms of these new functions, the initial condition is
\[\chi_i(0) = \chi_i(0) = \chi_i(0) = 0 \]  

[A-30]

[A-31]

Given the value of \(\chi_i\), the values of \(\chi_i(0)\) can be obtained by applying Eq. A-19. The value of \(\partial \chi_i(0)/\partial \xi^2\) can be approximated by taking finite differences
\[ \frac{d^2\chi_i(\Delta \xi, \tau)}{d\tau^2} = \frac{\chi_{i+1}(\tau) - 2\chi_i(\tau) + \chi_{i-1}(\tau)}{(\Delta \xi)^2} + O[(\Delta \xi)^3] \]  

where \( O[(\Delta \xi)^3] \) is a term of order \((\Delta \xi)^3\). If that term is neglected, Eq. A-16 becomes

\[ \frac{d\chi_i}{d\tau} = \left[ 1 + \frac{\rho}{(\chi_i + \rho)^2} \right]^{-1} \left( \chi_{i+1} - 2\chi_i + \chi_{i-1} \right) \]

for \( 0 < i < 2n \)  

[A-35]

For the boundary at \( i = 0 \), Eq. A-35 is slightly simplified. To accommodate Eq. A-24, one may define a fictitious \( \chi_{i-1} \) or \( \chi_{i+1} \). Then Eq. A-35 becomes

\[ \frac{d\chi_0}{d\tau} = \left[ 1 + \frac{\rho}{(\chi_0 + \rho)^2} \right]^{-1} \left( 2\chi_{i+1} - 2\chi_0 \right) \]

[A-36]

The most difficult part of the problem is implementing the boundary condition given by Eq. A-29. As a first step toward this, define \( s_i \) such that

\[ s_0 = s_{2n} = 1 \]

[A-37]

\[ s_i = 4 \] for \( i = 1, 3, ..., 2n - 1 \)

[A-38]

\[ s_i = 2 \] for \( i = 2, 4, ..., 2n - 2 \)

[A-39]

Then for any function \( f \), Simpson’s rule may be implemented as

\[ \int_0^h f(x) dx = \frac{\Delta x}{3} \sum_{i=0}^{2n} s_i f_i + O[(\Delta x)^3] \]  

[A-40]

where \( O[(\Delta x)^3] \) is a term of order \((\Delta x)^3\). If the higher-order term in Eq. A-40 is neglected, Eq. A-29 becomes, in discretized form

\[ \frac{d\chi_{2n}}{d\tau} = -\frac{\sigma \Delta \xi}{3\tau} \sum_{i=0}^{2n} s_i \frac{d\chi_i}{d\tau} \left[ 1 + \frac{\rho}{(\chi_i + \rho)^2} \right] \]

[A-41]

Note that \( d\chi_{2n}/d\tau \) appears on both sides of Eq. A-41. If Eq. A-41 is solved for \( d\chi_{2n}/d\tau \), one obtains

\[ \frac{d\chi_{2n}}{d\tau} = -\frac{\sigma \Delta \xi}{3\tau} \sum_{i=0}^{2n} s_i \frac{d\chi_i}{d\tau} \left[ 1 + \frac{\rho}{(\chi_i + \rho)^2} \right] \]

\times \sum_{i=0}^{2n-1} s_i \frac{d\chi_i}{d\tau} \left[ 1 + \frac{\rho}{(\chi_i + \rho)^2} \right] \]

[A-42]

Equations A-35, A-36, and A-42 provide a set of ordinary differential equations that may be solved to provide a discretized approximation to \( \chi_i(\xi, \tau) \).

A MathCAD program was written to solve these equations and to plot the predicted solution concentration of \( C_{x1} \) as a function of time. Fittings to the data have been made by changing the values of parameters \( D, \tau, \rho, \) and \( \eta \).

Asymptotic behavior (\( t \rightarrow \infty \)).—The concentration of \( C_{x1} \) in the solution at asymptotically large times can be computed by equating its concentration in the film at \( t = 0 \) to that in the film plus that in the solution at \( t \rightarrow \infty \). Thus,

\[ C_{x1} SL = C_{x1}(\infty) SL + C_{x1}(\infty) V \]

[A-43]

Note that \( C \) must become independent of \( \alpha \) as equilibrium is approached. In addition

\[ C(x, \infty) = C_x(x, \infty) + C_f(x, \infty) \]

[A-44]

and

\[ C_d(x, \infty) = C_d(\infty) \]

[A-45]

from which we obtain

\[ C_{x1} SL = [C_{x1}(\infty) + C_f(x, \infty)] SL + C_{x1}(\infty) V \]

[A-46]

We can eliminate \( C_a \) using Eq. A-7

\[ C(x, \infty) = \frac{\eta C_{x1}(\infty)}{C_{x1}(\infty) + \eta \rho} \]

[A-47]

Combining Eq. A-45 through A-47

\[ C_{x1} SL = C_{x1}(\infty)(V + SL) + \frac{\eta C_{x1}(\infty)}{C_{x1}(\infty) + \eta \rho} SL \]

[A-48]

Dividing through by \( SL \) and defining \( \sigma = V/SL \)

\[ C_{x1} = C_{x1}(\infty)(\sigma + 1) + \frac{\eta C_{x1}(\infty)}{C_{x1}(\infty) + \eta \rho} \]

[A-49]

Equation A-49 can be expressed as a quadratic expression in \( C_{x1}(\infty) \), i.e.

\[ C_{x1}(\infty) = C_{x1}(\infty)(\sigma + 1) + C_{x1}(\infty)(\sigma + 1) + \eta C_{x1}(\infty) \]

[A-50]

Putting Eq. A-50 into standard quadratic form

\[ (\sigma + 1)C_{x1}(\infty)^2 + [(\sigma + 1)\eta \rho + \eta - C_{x1}(\infty)] C_{x1}(\infty) - C_{x1}(\infty) = 0 \]

[A-51]

Equation A-51 has two roots, only one of which is physically acceptable, the other being negative.

References